

Assistance of Core-Ion Dissociation by Energy Transmission from Solvent Molecules in $\text{N}_3\text{O}_3^-(\text{NO})_n$ ($n = 0\text{--}38$) Impact onto Silicon Surface

Hitoshi Yamaguchi,[#] Hisato Yasumatsu,[†] and Tamotsu Kondow^{*†}

East Tokyo Laboratory, Genesis Research Institute, Inc., 717-86 Futamata, Ichikawa, Chiba 272-0001

[†]Cluster Research Laboratory, Toyota Technological Institute: in East Tokyo Laboratory, Genesis Research Institute, Inc., 717-86 Futamata, Ichikawa, Chiba 272-0001

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A size-selected cluster anion, $\text{N}_3\text{O}_3^-(\text{NO})_n$ ($n = 0\text{--}38$), was allowed to collide onto a silicon surface, and the product anions, NO_2^- , N_2O_2^- and $\text{N}_3\text{O}_3^-(\text{NO})_m$ ($m = 0\text{--}2$), from the surface were observed mass-spectroscopically. The branching fractions of the product anions were measured as functions of the energy of the incoming parent cluster anion and the size, n . It was found that NO molecules located in the vicinity of the core ion, N_3O_3^- , transmit their kinetic energy to N_3O_3^- most efficiently.

In the collision of a cluster ion onto a solid surface, the kinetic energy of the incoming cluster ion is converted very efficiently to its internal energy,¹⁻³ especially to that of its core ion. In a surface impact of $\text{I}_2^-(\text{CO}_2)_n$, the kinetic energy of the solvent CO_2 molecules at the waist positions of the core ion, I_2^- , is transmitted intensively to the I_2^- bond, and as a result, the bond splits as if a piece of wood splits by a hammer thrust against a wedge.⁴⁻⁶ This finding shows that the energy transmission depends critically on the solvation structure and the number of the solvent molecules.⁴⁻⁷ It is important to know how much energy is transmitted to the core ion from a solvent molecule of a known location in the solvation structure. This energy transmission is appropriately observed in the collision of $\text{N}_3\text{O}_3^-(\text{NO})_n$ onto a solid surface through monitoring the branching fractions of the two dissociation channels of the core ion, N_3O_3^- , into NO_2^- and N_2O_2^- .⁸⁻¹⁰ Note that the cluster anion consists of a 'hard' core ion, N_3O_3^- , and a 'soft' solvation shell, $(\text{NO})_n$.^{11,12} In this connection, we studied the impact of size-selected $\text{N}_3\text{O}_3^-(\text{NO})_n$ onto a silicon surface.

The measurements were performed in a tandem time-of-flight (TOF) mass spectrometer equipped with a cluster ion source and an ultrahigh-vacuum surface-collision chamber.⁴⁻⁶ A gas mixture of 5–20% nitric oxide (Nihon Sanso, 97% pure) in argon (Nihon Sanso, 99.9995% pure) having a pressure of 4 atm was passed through a coiled tube cooled in a slush bath of an ethanol/liquid-nitrogen mixture at a temperature of ~ 170 K; impurities such as nitrogen dioxide and water were eliminated. The gas mixture was expanded supersonically into the source chamber from a pulsed valve. Cluster anions, $\text{N}_3\text{O}_3^-(\text{NO})_n$, were produced by injecting electrons (50–300 eV) into the supersonic expansion region,¹³ and a cluster ion of a desired size was selected by using a mass gate placed in the primary TOF mass spectrometer. The cluster anion, $\text{N}_3\text{O}_3^-(\text{NO})_n$, with a given n ($n = 0\text{--}38$) was allowed to collide onto a silicon surface at a collision energy, E_{col} , per core ion. The uncertainty in E_{col} was determined by a retarding potential method for every cluster size; typically, 1.8 eV (one-standard deviation) in the collision of $\text{N}_3\text{O}_3^-(\text{NO})_8$. The product anions from the surface were mass-analyzed in the secondary TOF mass spectrom-

eter. It was found that the detector collects more than 90% of the product anions. All the measurements were performed at a temperature of 300 K and a pressure less than 4×10^{-8} Pa. The silicon surface (25×15 mm²) was prepared by heating a p-doped Si(100) wafer (specific resistance of 4.5–6 Ω cm) covered with silicon oxide layers (9-nm thickness) at 900 K for 10 hours under an ambient pressure of $\sim 10^{-7}$ Pa or less.

Upon the impact of $\text{N}_3\text{O}_3^-(\text{NO})_n$, NO_2^- , N_2O_2^- and $\text{N}_3\text{O}_3^-(\text{NO})_m$ ($m = 0, 1, 2$), together with a trace amount of NO^- and O^- , were scattered from the surface. The branching fractions, f_x ($x = 1, 2, 3$), of the product anions, NO_2^- , N_2O_2^- and $\text{N}_3\text{O}_3^-(\text{NO})_m$, respectively, are given by

$$f_x = \frac{[\text{Y}]}{[\text{NO}_2^-] + [\text{N}_2\text{O}_2^-] + \sum_{m=0}^2 [\text{N}_3\text{O}_3^-(\text{NO})_m]}, \quad (1)$$

where [Y] represents the intensity of a given product anion, Y.

The branching fractions, f_x ($x = 1, 2, 3$), for the N_3O_3^- impact are plotted against the collision energy, E_{col} , per N_3O_3^- in Figure 1. The branching fraction (f_1) for the NO_2^- production starts to increase in the vicinity of the smallest collision energy studied (~ 6 eV); the appearance energy of NO_2^- is determined to be 6 ± 8 eV. Note that NO_2^- is produced from N_3O_3^- specifically in the surface impact probably because of impulsive energy transmission to the internal modes of N_3O_3^- . On the other hand, the

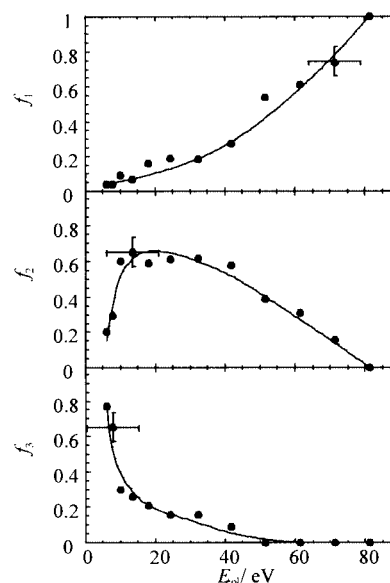


Figure 1. Branching fractions, f_1 , f_2 and f_3 , for the production of NO_2^- , N_2O_2^- and N_3O_3^- , respectively, as a function of the collision energy, E_{col} , in the impact of N_3O_3^- . The solid curves represent eye guides. The definition of the branching fractions is given in text.

branching fraction (f_2) for the N_2O_2^- production rises sharply at the smallest collision energy studied. This finding is consistent with the barrier height of less than 2.5 eV⁸ for the reaction,

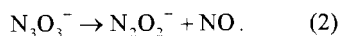


Figure 2 shows f_1 and f_2 as a function of the collision energy, E_{col} , when $\text{N}_3\text{O}_3^-(\text{NO})_n$ is allowed to collide onto the silicon surface. A reaction mechanism similar to that for the N_3O_3^- impact is operative because $\text{N}_3\text{O}_3^-(\text{NO})_n$ consists of a hard core ion, N_3O_3^- , and a soft solvation shell, $(\text{NO})_n$,^{11,12} that is, N_3O_3^- dissociates by the surface impact with the assistance of kinetic-energy transmission from NO molecules in the shell. Actually, a less collision energy is needed to obtain the same f_1 and f_2 values if the core ion possesses a more number of NO molecules (see Figure 2). In other words, the core ion, N_3O_3^- , gains an effective collision energy, E_{eff} ; the sum of prompt energy gain by surface collision of the parent cluster anion and subsequent energy gain, E_{solv} , through intracluster collision with the solvent NO molecules. The prompt energy gain is given by the collision energy, E_{col} . The effective collision energy is then expressed as

$$E_{\text{eff}} = E_{\text{col}} + E_{\text{solv}}. \quad (3)$$

As the same reaction mechanism is operative, irrespective of the number of the solvent NO molecules as described above, the dependences of f_1 and f_2 on E_{eff} must be expressed by a universal relationship between f_x ($x = 1, 2$) and E_{eff} valid at any number of the solvent NO molecules.

Let us define the efficiency of the energy transmission as

$$\eta = E_{\text{solv}} / E_{\text{col}}. \quad (4)$$

Then, E_{eff} is given from equations (3) and (4) as

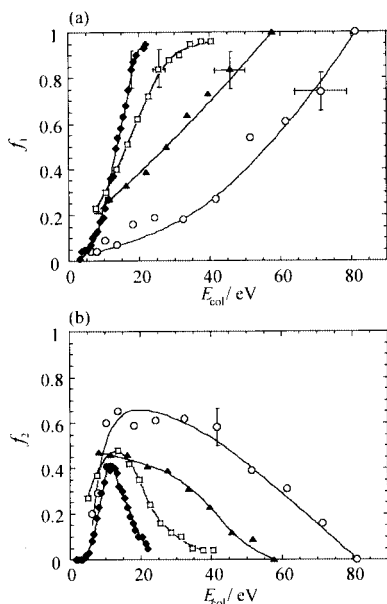


Figure 2. Branching fractions, f_1 and f_2 , for the production of NO_2^- (panel a) and N_2O_2^- (panel b) from $\text{N}_3\text{O}_3^-(\text{NO})_n$ as a function of the collision energy, E_{col} , (per core ion, N_3O_3^-). The marks, \circ , \blacktriangle , \square and \blacklozenge show the branching fractions for $n = 0, 2, 7$ and 38 , respectively. The solid curves represent eye guides. The definition of the branching fractions is given in text.

$$E_{\text{eff}} = E_{\text{col}}(\eta + 1). \quad (5)$$

The efficiency, η , is determined so that the dependences of f_1 and f_2 on E_{eff} are expressed by the universal relationship. As η is zero for the N_3O_3^- impact or for $n = 0$, the universal relationship coincides with that for the N_3O_3^- impact. Figure 3 shows the η value plotted against the number of the solvent NO molecules. The η values obtained from the f_1-E_{col} and f_2-E_{col} relations coincide. The steep slope in the $n < \sim 8$ range and the gentle slope in the $n > \sim 10$ range indicate that NO molecules in the vicinity of the core ion transmit their kinetic energy much more efficiently than those apart from the core ion. The slope of the $\eta-n$ plot in the $n < \sim 8$ range (see Figure 3) shows that the effective collision energy of the core ion (E_{eff}) increases by 25% per one NO molecule located in the close vicinity of the core ion. The f_2-E_{col} relation at energies lower than ~ 10 eV was ruled out because NO release from a parent cluster anion dominates in this energy range.

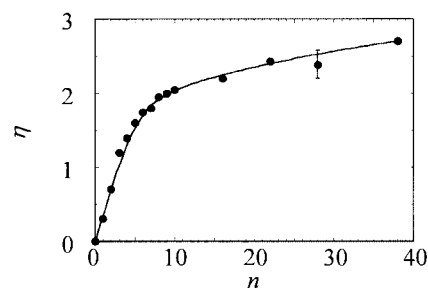


Figure 3. Efficiency, η , for energy transmission from the solvent NO molecules to the core ion, N_3O_3^- , as a function of the number of the solvent NO molecules, n . The solid curve represents an eye guide. The definition of η is given in text.

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References and Notes

- # Present address: Bridgestone Corporation, 3-1-1 Ogawahigashi-cho, Kodaira, Tokyo 187-8531.
- 1 C. L. Cleveland and U. Landman, *Science*, **257**, 355 (1992).
- 2 I. Schek and J. Jortner, *J. Chem. Phys.*, **104**, 4337 (1996).
- 3 T. Raz and R. D. Levine, *Chem. Phys.*, **213**, 263 (1996).
- 4 H. Yasumatsu, A. Terasaki, and T. Kondow, *J. Chem. Phys.*, **106**, 3806 (1997).
- 5 H. Yasumatsu, U. Kalmbach, S. Koizumi, A. Terasaki, and T. Kondow, *Z. Phys. D*, **40**, 51 (1997).
- 6 U. Kalmbach, H. Yasumatsu, S. Koizumi, A. Terasaki, and T. Kondow, *J. Chem. Phys.*, **110**, 7443 (1999).
- 7 I. Schek, J. Jortner, T. Raz, and R. D. Levine, *Chem. Phys. Lett.*, **257**, 273 (1996).
- 8 A. Snis and I. Panas, *Chem. Phys. Lett.*, **305**, 285 (1999).
- 9 C. L. Lugez, W. E. Thompson, M. E. Jacox, A. Snis, and I. Panas, *J. Chem. Phys.*, **110**, 10345 (1999).
- 10 T. Tsukuda and T. Nagata, unpublished data.
- 11 T. Tsukuda, M. Saeki, L. Zhu, and T. Nagata, *Chem. Phys. Lett.*, **295**, 416 (1998).
- 12 H. S. Carman, Jr, *J. Chem. Phys.*, **100**, 2629 (1994).
- 13 M. A. Johnson and W. C. Lineberger, in "Techniques of Chemistry," ed. by J. M. Farrar and W. H. Saunders, Wiley, New York (1988), Vol. 20, p 591.